



IMEDEA



POLARIZATION COUPLING AND TRANSVERSE PATTERNS IN TYPE-II OPTICAL PARAMETRIC OSCILLATORS

Gonzalo Izús

Maxi San Miguel

Daniel Walgraef

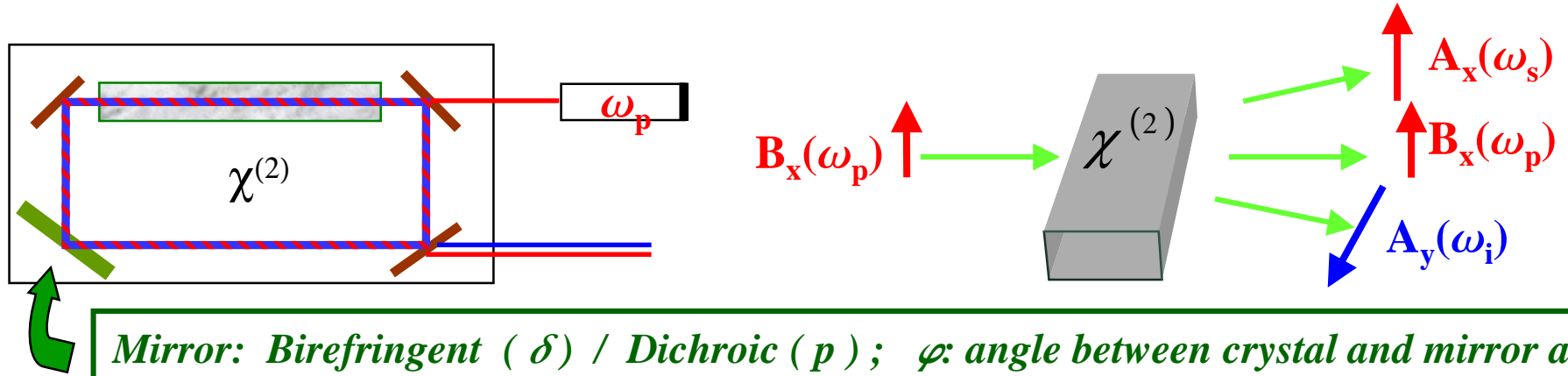
Marco Santagiustina

Opt. Lett. 25, 1454 (2000);

Phys. Rev. E, 64, 056231-1/15 (2001);

Phys. Rev. E (2002)

Polarization Coupling in Type-II OPO



$$\begin{aligned} \partial_t B_x &= \gamma'_x \left[- (1 + i\Delta'_x) B_x + E_0 + ia'_x \nabla^2 B_x + 2iK_0 A_x A_y + c'_x B_y \right] \\ \partial_t B_y &= \gamma'_y \left[- (1 + i\Delta'_y) B_y + ia'_y \nabla^2 B_y + c'_y B_x \right] \\ \partial_t A_x &= \gamma_x \left[- (1 + i\Delta_x) A_x + ia_x \nabla^2 A_x + iK_0 A_y^* B_x + c_x A_y \right] \\ \partial_t A_y &= \gamma_y \left[- (1 + i\Delta_y) A_y + ia_y \nabla^2 A_y + iK_0 A_x^* B_x + c_y A_x \right] \end{aligned}$$

$$c_{x,y} = \frac{(p + i\delta) \sin(2\varphi)}{T \pm p \cos(2\varphi)}$$

$$c_x = -c_y^* = \epsilon_0 e^{i\theta}$$

Intracavity $\lambda/4$ plate

$c = 0$: $A_x \rightarrow A_x e^{i\varphi}$, $A_y \rightarrow A_y e^{-i\varphi} \longrightarrow$ No relative phase preferred

$$\Delta_e = (\gamma_x \Delta_x + \gamma_y \Delta_y) / (\gamma_x + \gamma_y) \begin{cases} \Delta_e > 0 : \text{Homogeneous solutions selected} \\ \Delta_e < 0 : \text{Phase patterns (Travelling Waves)} \end{cases}$$

$c \neq 0$: **Phase locked solutions.** E. Manson, and N. Wong Opt. Lett. **23**, 1733 (1998)
C. Fabre, E. Manson, and N. Wong Opt. Comm. **170**, 299 (1999)

Transverse effects? : Phase Walls, Threshold lowering, Standing Waves

Type-II Optical Parametric Oscillators: homogeneous solutions ($\Delta_e > 0$)

A) No linear polarization coupling: $c = 0$

• Threshold of instability of $A_x = A_y = 0$ \longrightarrow

$$|F_c|^2 = |E_0 K_0 / (1 + i \Delta_x)|^2 = 1 + \Delta_e^2$$

$$q_0 = 0$$

$$\omega_0 = \gamma_x \gamma_y (\Delta_x - \Delta_y) / (\gamma_x + \gamma_y)$$

B) Linear polarization coupling: $c \neq 0$ ($\Delta_{x,y} > 0$)

$$c_x = c_y = c_r + i c_i \quad (\varphi = 45^\circ)$$

Stationary uniform phase locked solution

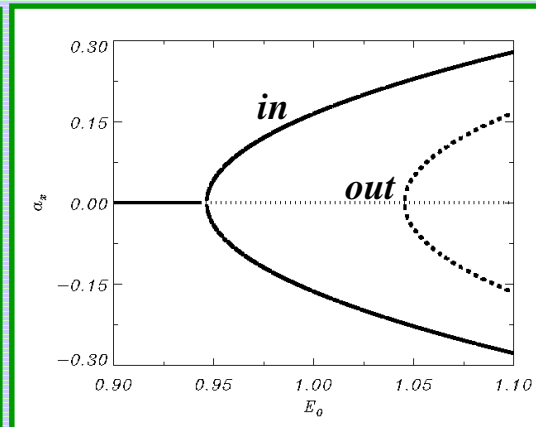
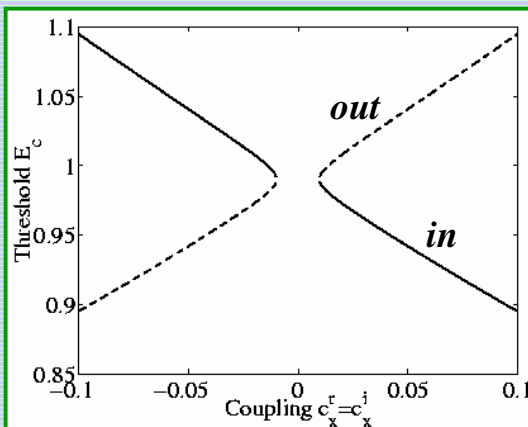
$$A_x = a_x \exp(i \psi_x), \quad A_y = a_y \exp(i \psi_y)$$

- $|E| > E_c$
- $4 (c_r + c_i \Delta_x) (c_r + c_i \Delta_y) \geq (\Delta_y - \Delta_x)^2$

$$\bullet \sin(\psi_y - \psi_x) = (\Delta_x - \Delta_y) / [4 (c_r + c_i \Delta_x) (c_r + c_i \Delta_y)]^{1/2}$$

$$\bullet a_y^2 = \Gamma a_x^2 \quad \Gamma = (c_r + c_i \Delta_x) / (c_r + c_i \Delta_y)$$

THRESHOLD



- Well within the locked regime: “In” ($\psi_y - \psi_x \cong 0$) and “Out” ($\psi_y - \psi_x \cong \pi$) phase solutions of different threshold
- Two equivalent “in” homogeneous solutions (+/-) $A_{x,y}^+ = -A_{x,y}^-$
- Domain walls between A^+ y A^-

Polarization of Phase Locked States and Domain Walls

A) Homogeneous phase locked solutions

Polarization state determined by locked value of $\psi_y - \psi_x$

Stokes parameters

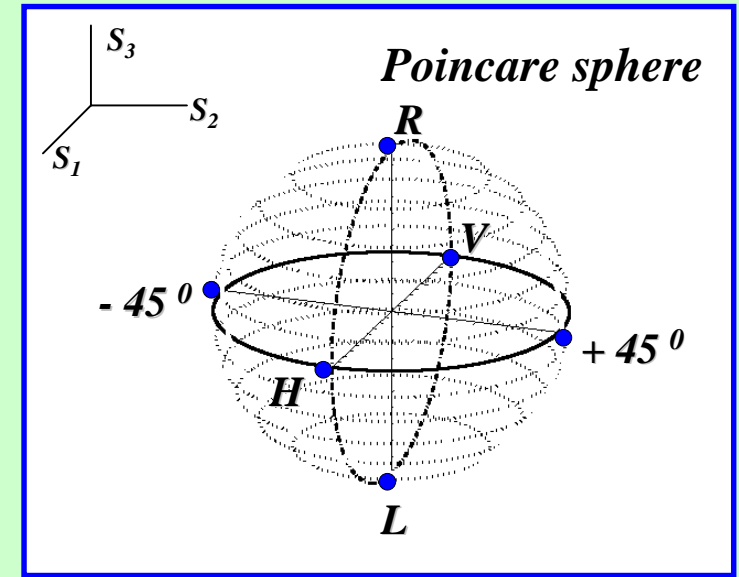
$$S_1 = (1 - \Gamma) / (1 + \Gamma)$$

$$S_2 = [2\Gamma^{1/2} / (1 + \Gamma)] \cos(\psi_y - \psi_x)$$

$$S_3 = [-2\Gamma^{1/2} / (1 + \Gamma)] \sin(\psi_y - \psi_x)$$

$$\Gamma = (c_r + c_i \Delta_x) / (c_r + c_i \Delta_y)$$

$$a_y^2 = \Gamma a_x^2$$



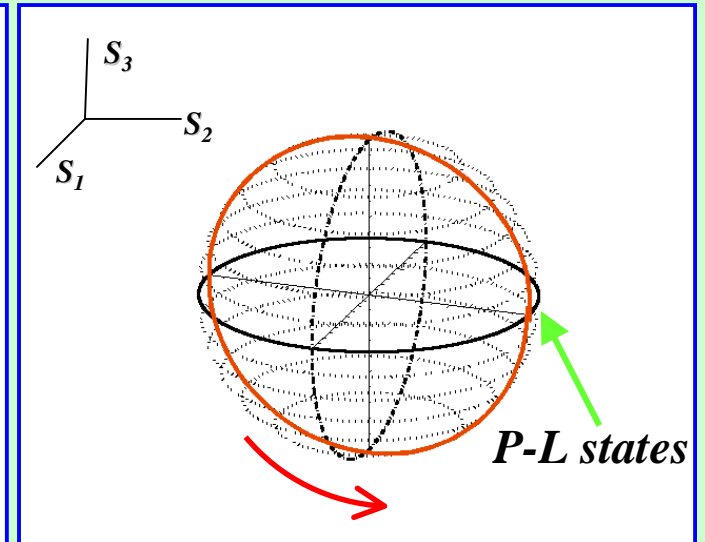
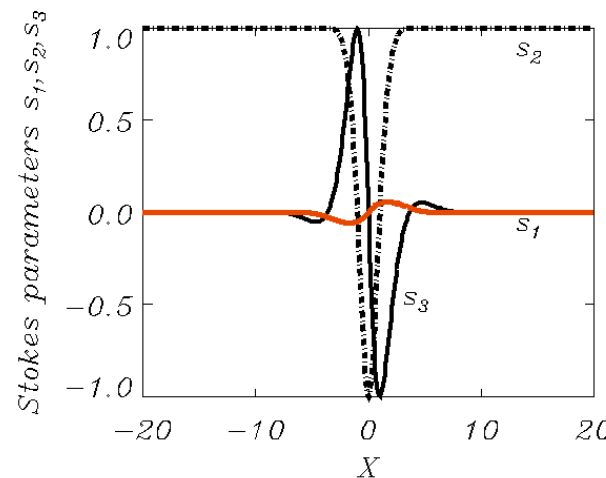
B) Phase polarization Bloch domain walls

Deeply in the P-L regime

$$\Delta_x = \Delta_y, \Gamma = 1, \psi_y = \psi_x$$

$$(S_1, S_2, S_3) = (0, 1, 0)$$

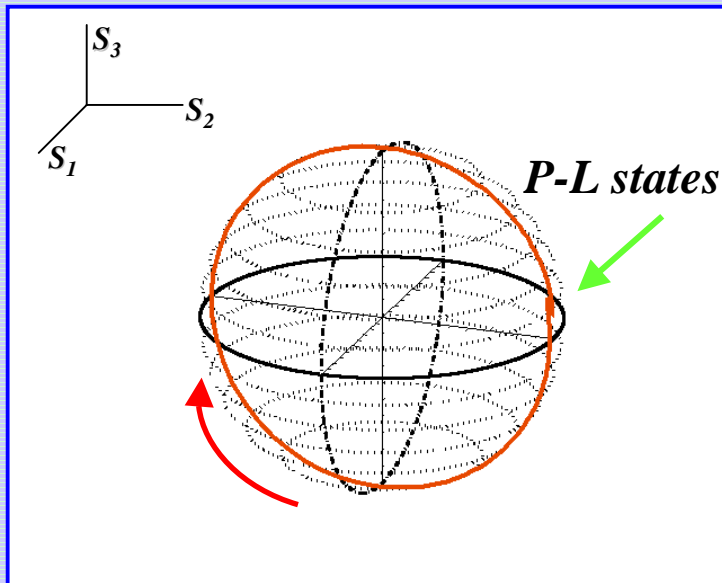
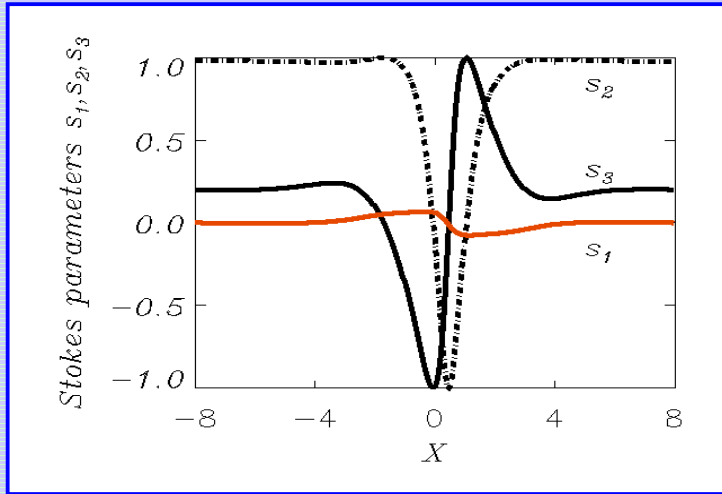
LPS at: $\theta = \pi/4$ ($-3\pi/4$)



Polarization of Domain Walls

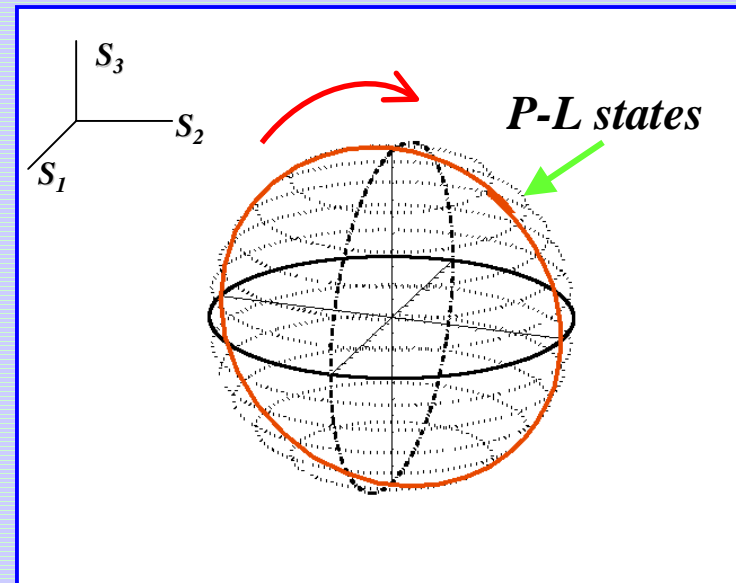
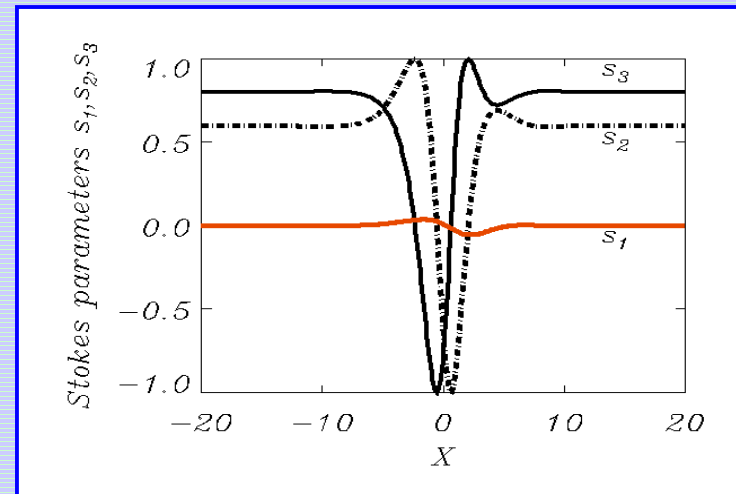
• Phase Locked state

$$\Delta_x \neq \Delta_y, \Gamma = 1, \quad \sin(\psi_y - \psi_x) = 0.2$$



• Near the P-L transition

$$\Delta_x \neq \Delta_y, \Gamma = 1, \quad \sin(\psi_y - \psi_x) = 0.8$$

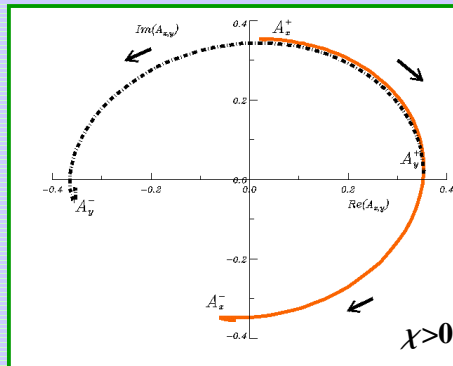
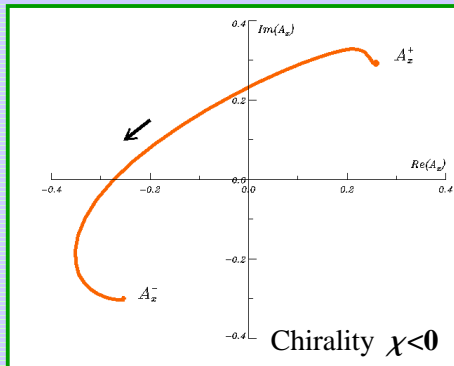
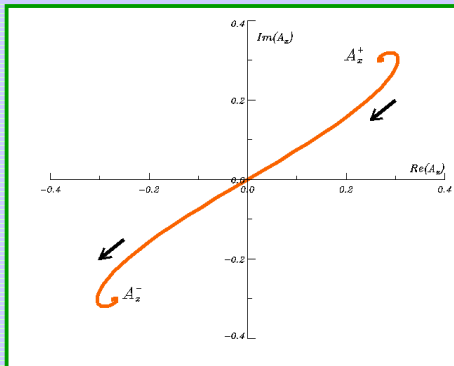
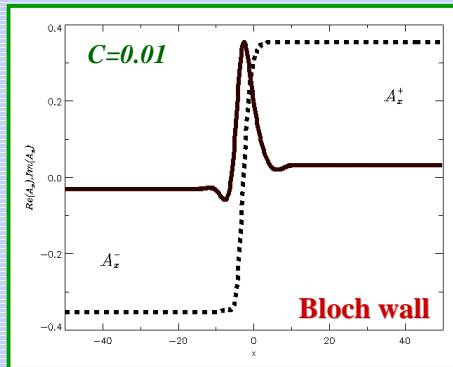
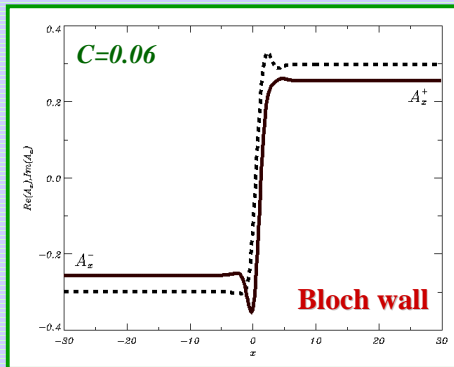
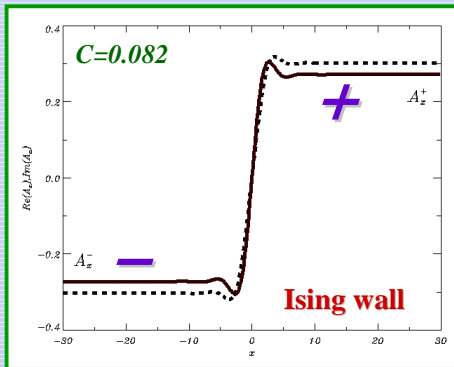


d=1 Domain Walls

G. Izús, M. Santagiustina, and M. San Miguel, *Opt. Lett.* **23**, 1167 (2000)

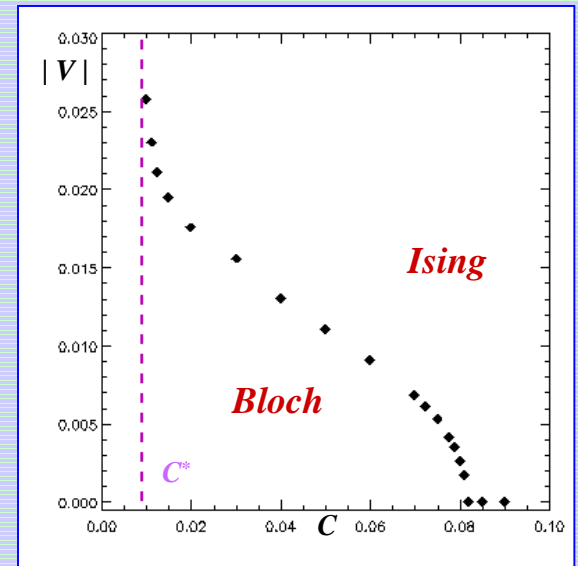
G. Izús, M. San Miguel, and M. Santagiustina, *submitted* (2001)

Bloch-Ising transition controlled by polarization coupling



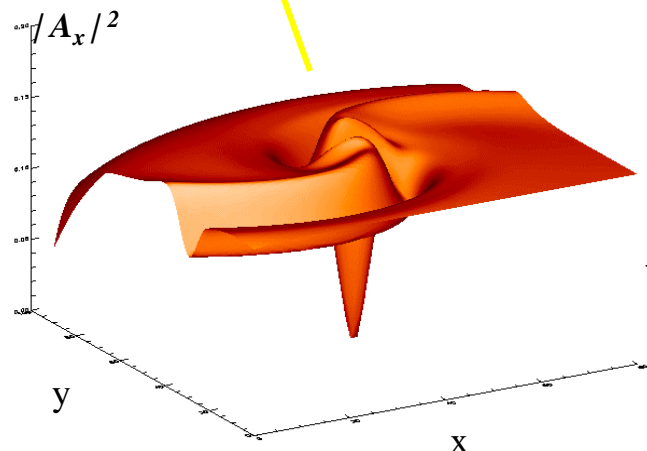
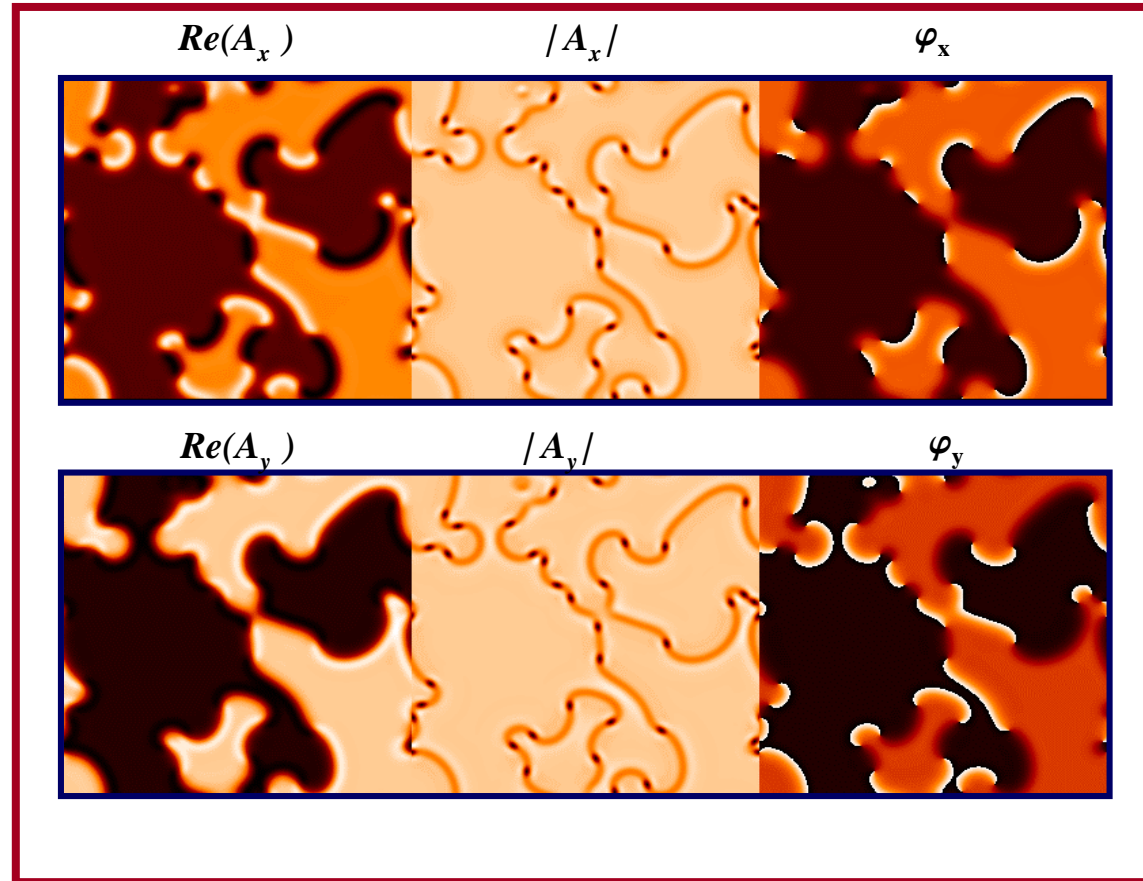
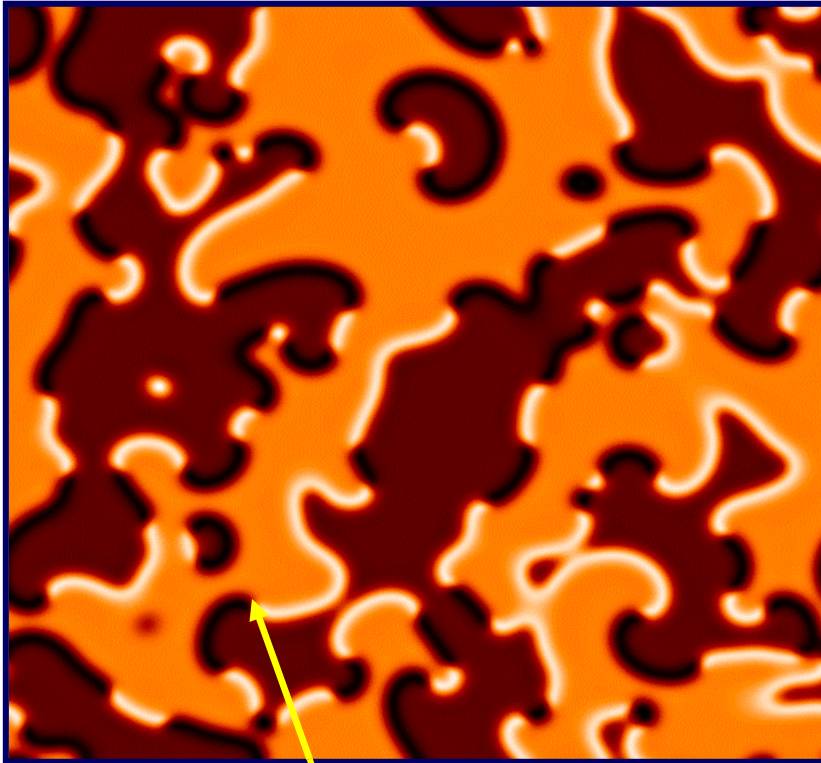
Two dynamical regimes

- $\gamma_x \Delta_x = \gamma_y \Delta_y$ static Bloch wall
- $\gamma_x \Delta_x \neq \gamma_y \Delta_y$ Bloch walls move in opposite directions for opposite χ



$d=2$ BLOCH WALLS

$Re(A_x)$

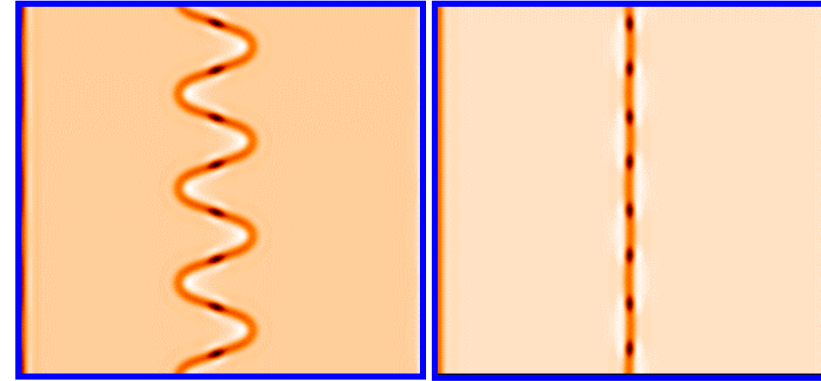


Defect: Point of change of chirality on the wall

BLOCH WALL DYNAMICS (I)

$$\gamma_x \Delta_x = \gamma_y \Delta_y \text{ Flat wall stable}$$

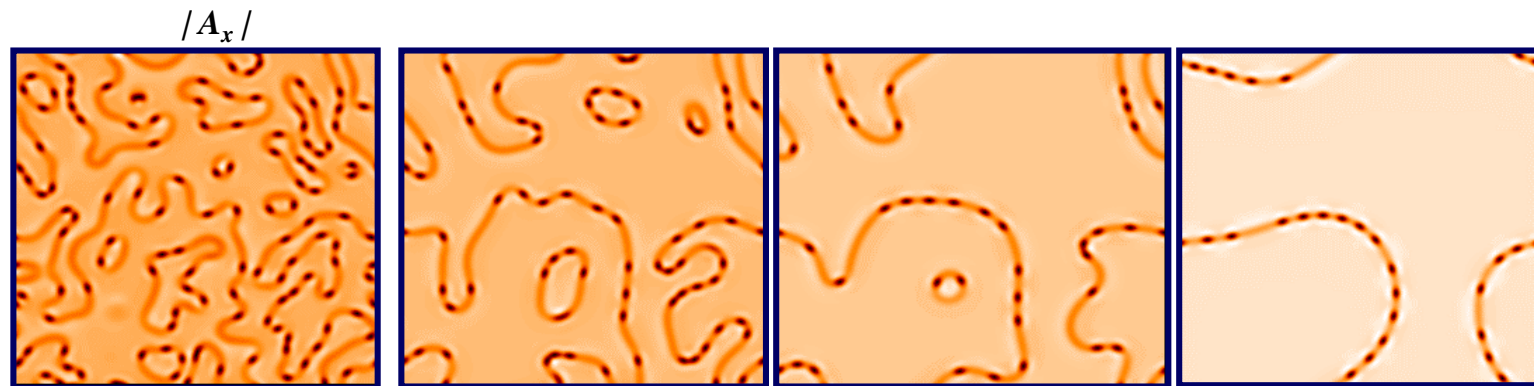
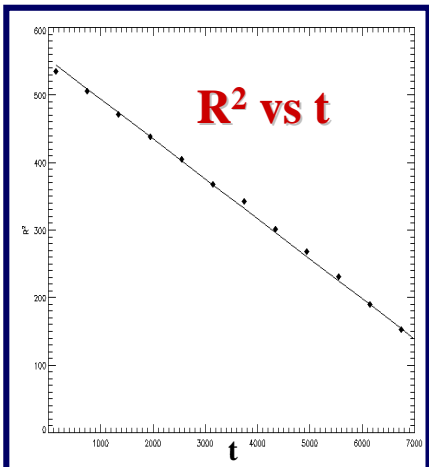
$|A_x|$



t=0

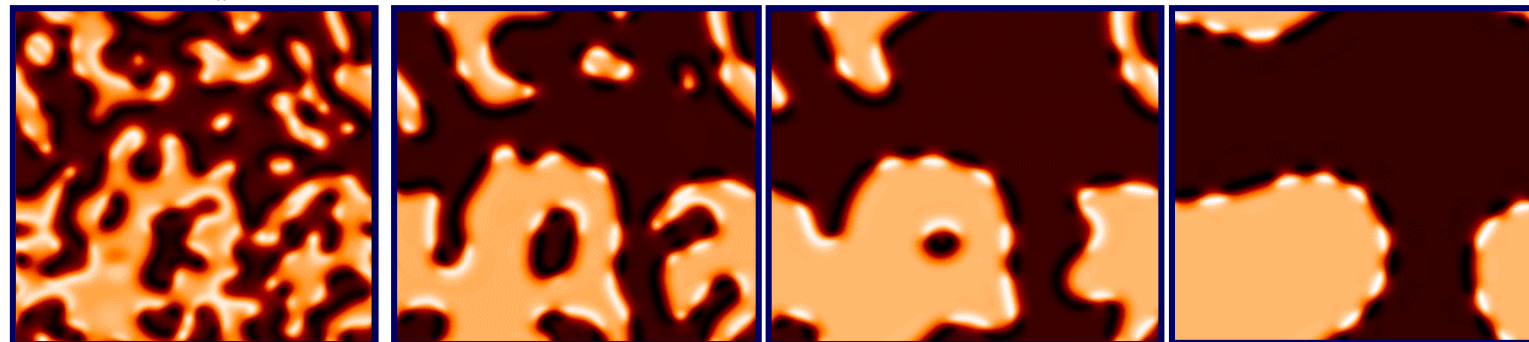
t=2000

*Domain Growth
controlled by
curvature*



$Re(A_x)$

time



t=100

t=350

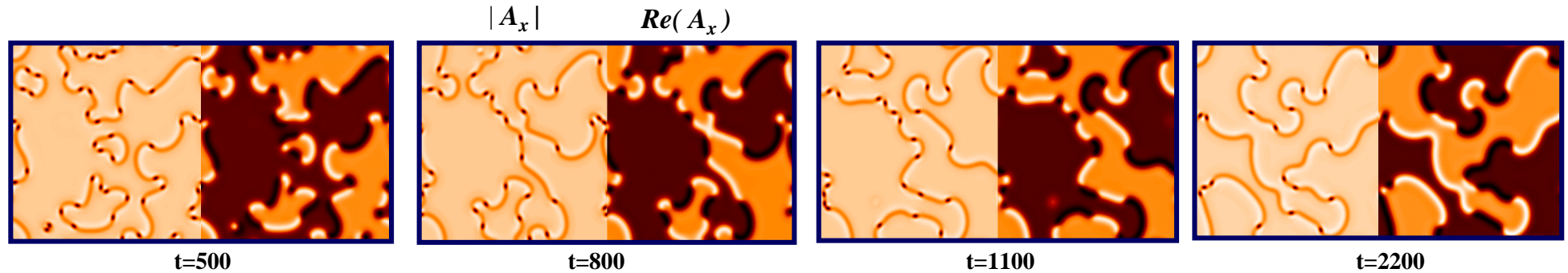
t=750

t=2900

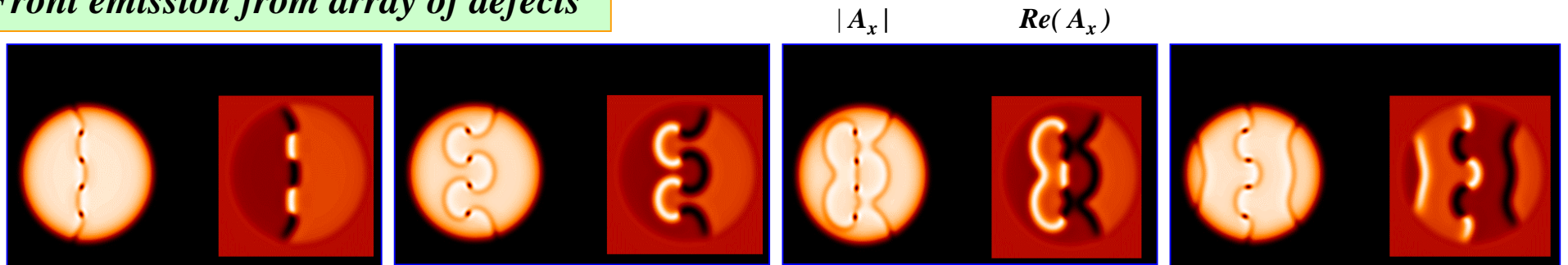
BLOCH WALL DYNAMICS (II)

$\gamma_x \Delta_x \neq \gamma_y \Delta_y$ Walls with different
chirality move in opposite directions

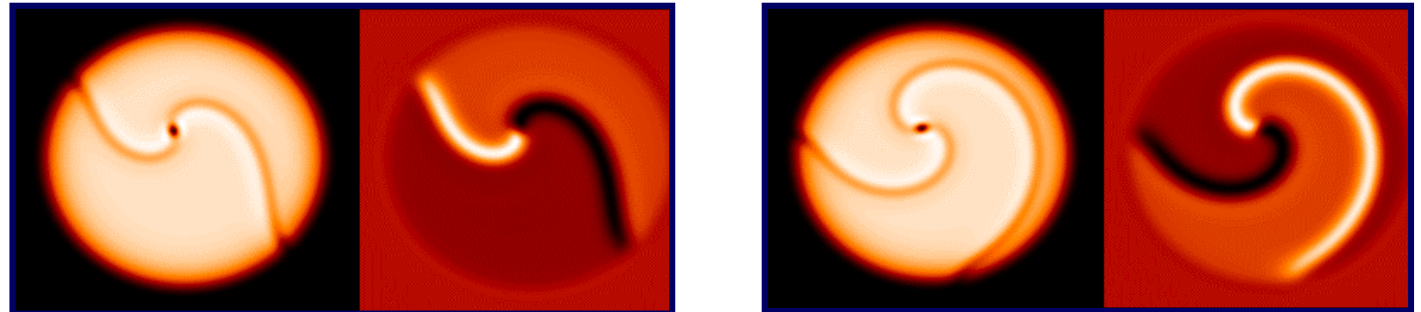
Persistent Dynamics



Front emission from array of defects

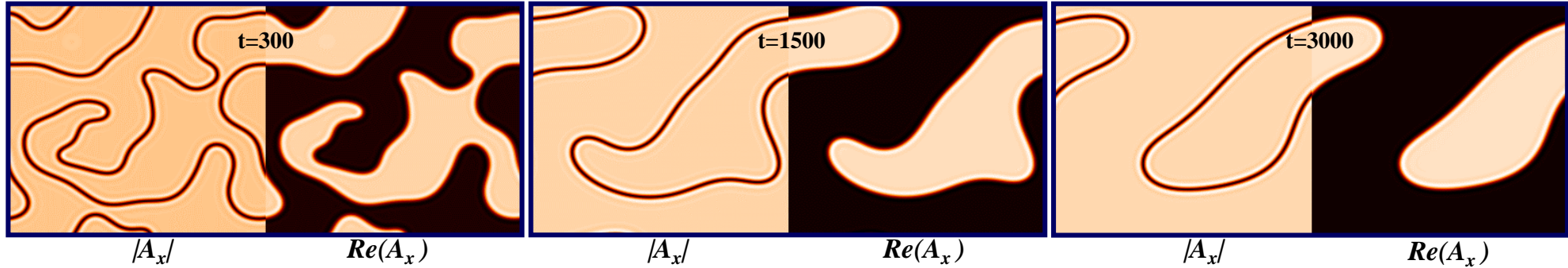


Two armed
rotating spiral
centered in defect

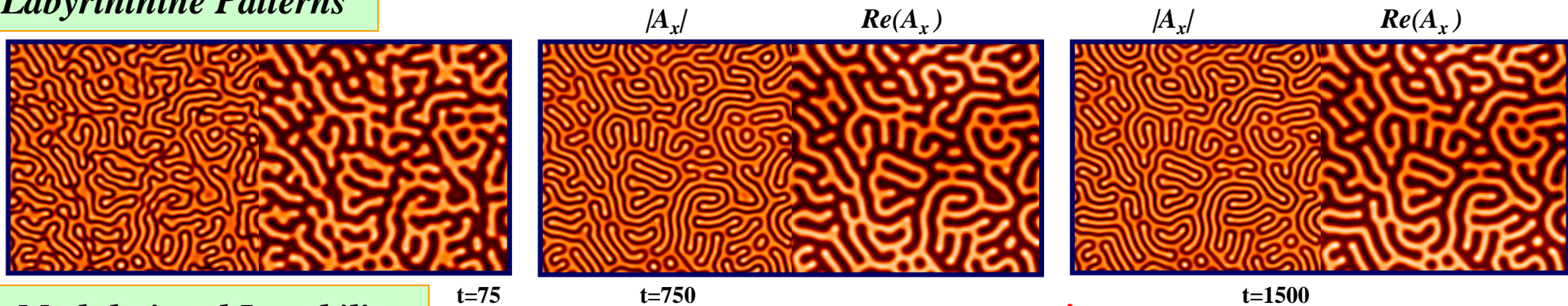


ISING WALL DYNAMICS

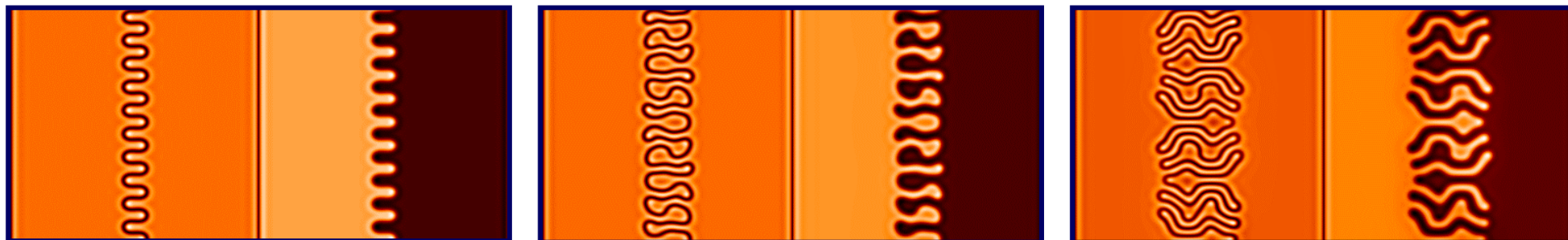
Coarsening near the Bloch-Ising transition



Labyrinthine Patterns



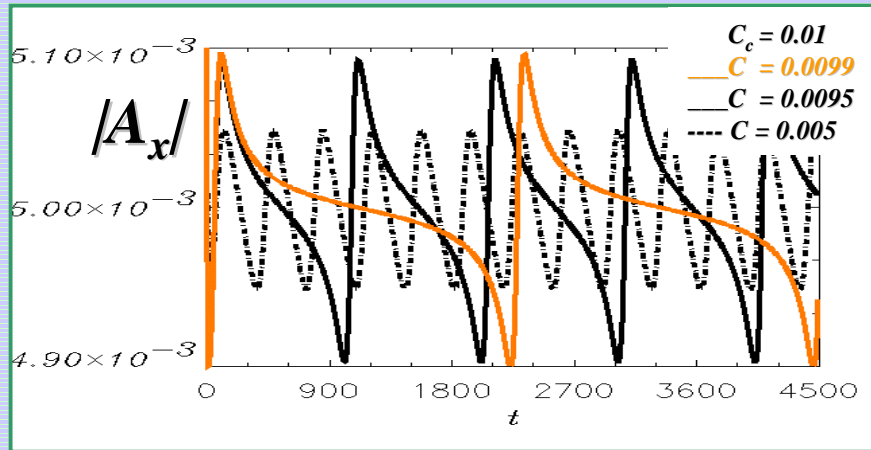
Modulational Instability



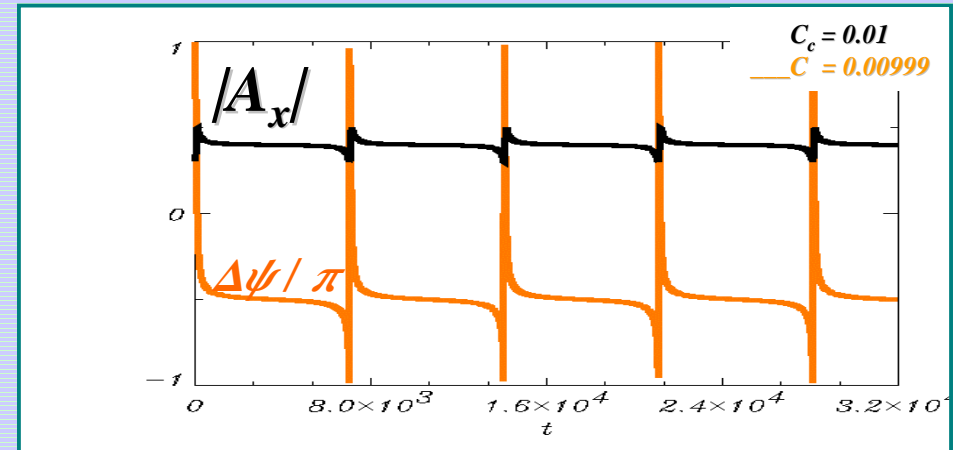
Time Oscillatory Bloch Domain Walls

Outside the phase-locked regime

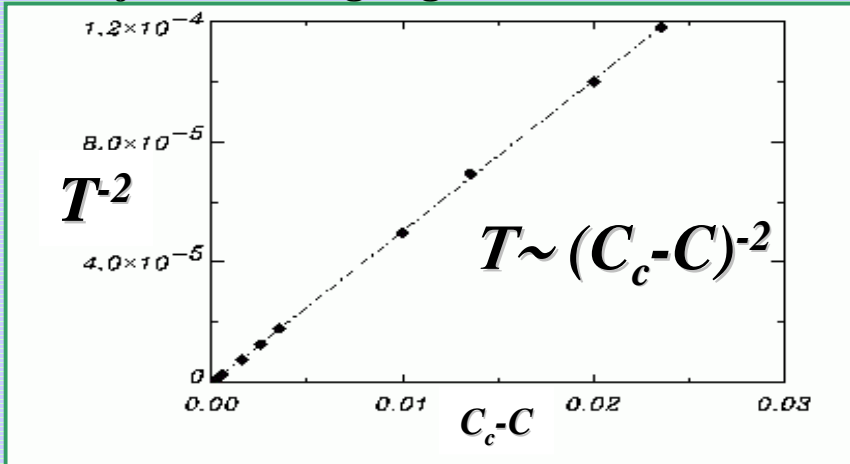
• Time oscillatory uniform state



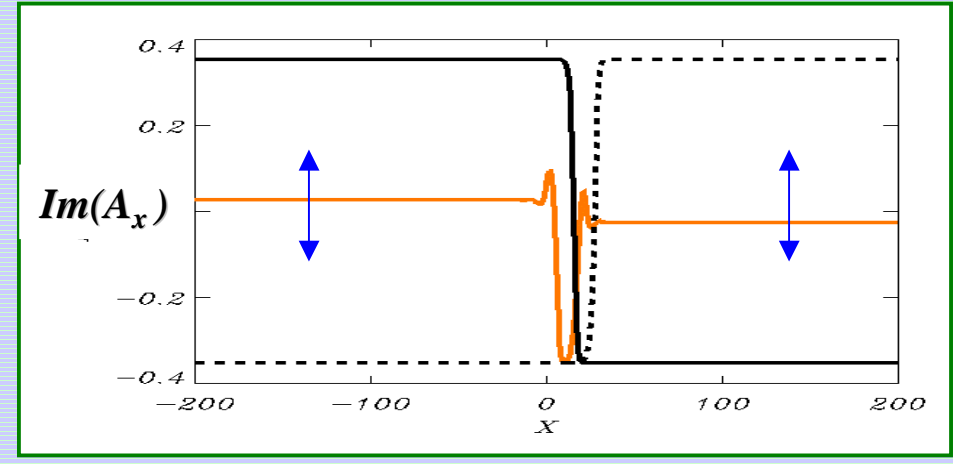
• Relative phase $\Delta\psi$ oscillates in time



• The period of time oscillation diverges in the limit of the locking regime



• Oscillatory Bloch domain walls: **Dynamics** ?



Type-II Optical Parametric Oscillators: Pattern Formation ($\Delta_e < 0$)

A) No linear polarization coupling: $c = 0$

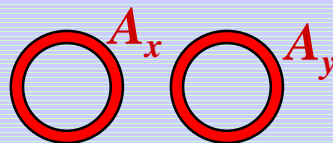
- Threshold of instability of $A_x = A_y = 0$

$$|F_c|^2 = 1, \quad q_0^2 = -\Delta_e / (\gamma_x a_x + \gamma_y a_y),$$

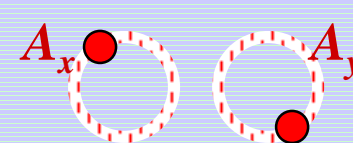
$$\omega_0 = \gamma_x \gamma_y [\Delta_x - \Delta_y + q_0^2 (a_x - a_y)] / (\gamma_x + \gamma_y)$$

Far Field Solutions

Below Threshold



Above Threshold

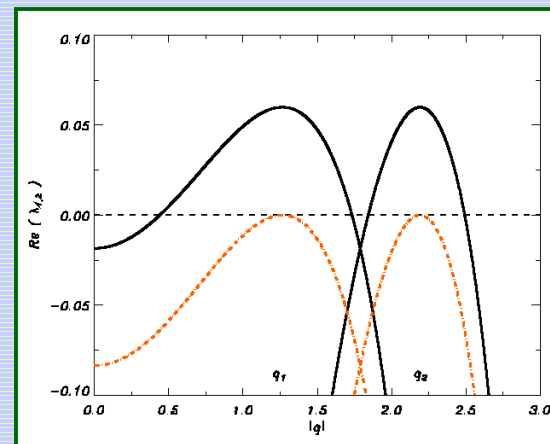
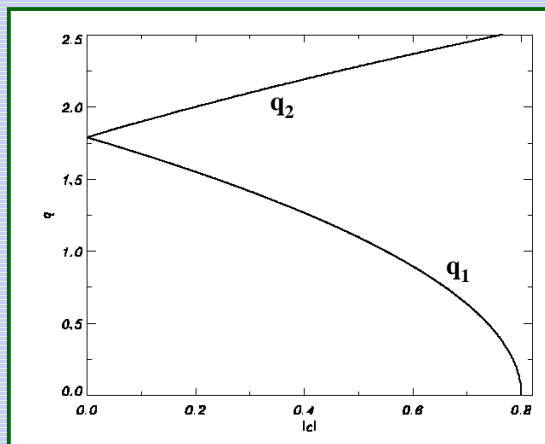
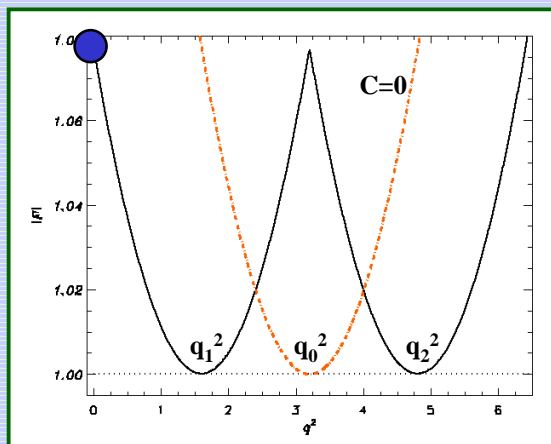


B) Linear polarization coupling $c_x = -c_y^* = c$ (real)

THRESHOLD for symmetric coefficients $\Delta_x = \Delta_y, \gamma_x = \gamma_y, a_x = a_y$

$$|F_c|^2 = 1, \quad \omega_0 = 0$$

$$q^2_{1,2} = (-\Delta \pm c) / a$$



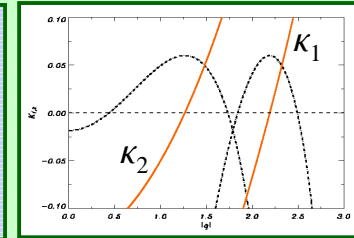
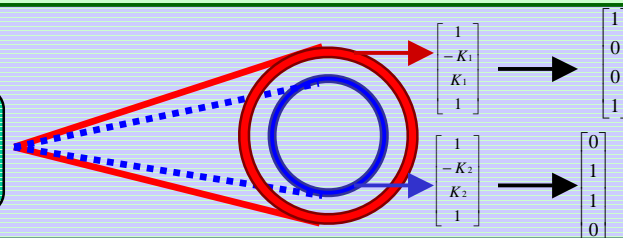
- Lower threshold for pattern formation: two competing modes with equal growth rate
- Uniform phase locked solutions ($q=0$) (Fabre et. al. Opt. Comm. 170, 299 (1999))

Type-II OPO, $\Delta_e < 0$: Symmetric coefficients

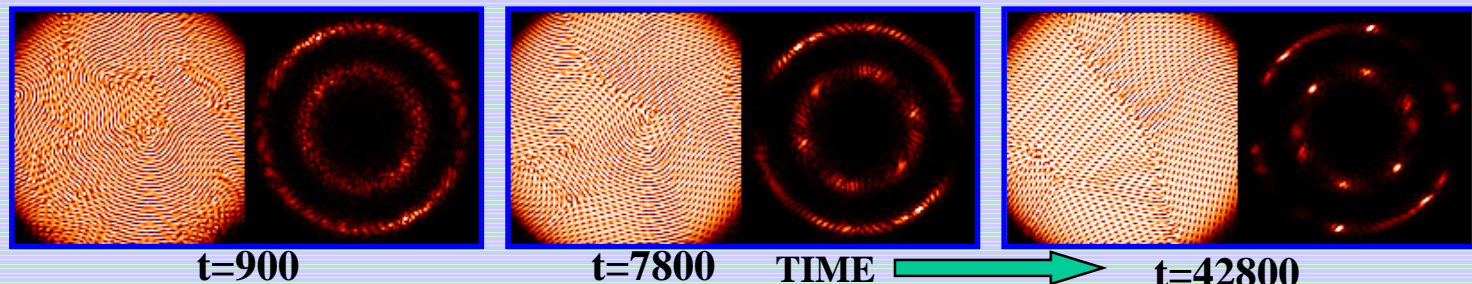
Linear stability analysis

$$\frac{\partial}{\partial t} \begin{bmatrix} \Re(A_x) \\ \Im(A_x) \\ \Re(A_y) \\ \Im(A_y) \end{bmatrix} = \hat{L}(c) \begin{bmatrix} \Re(A_x) \\ \Im(A_x) \\ \Re(A_y) \\ \Im(A_y) \end{bmatrix}$$

OPO-II

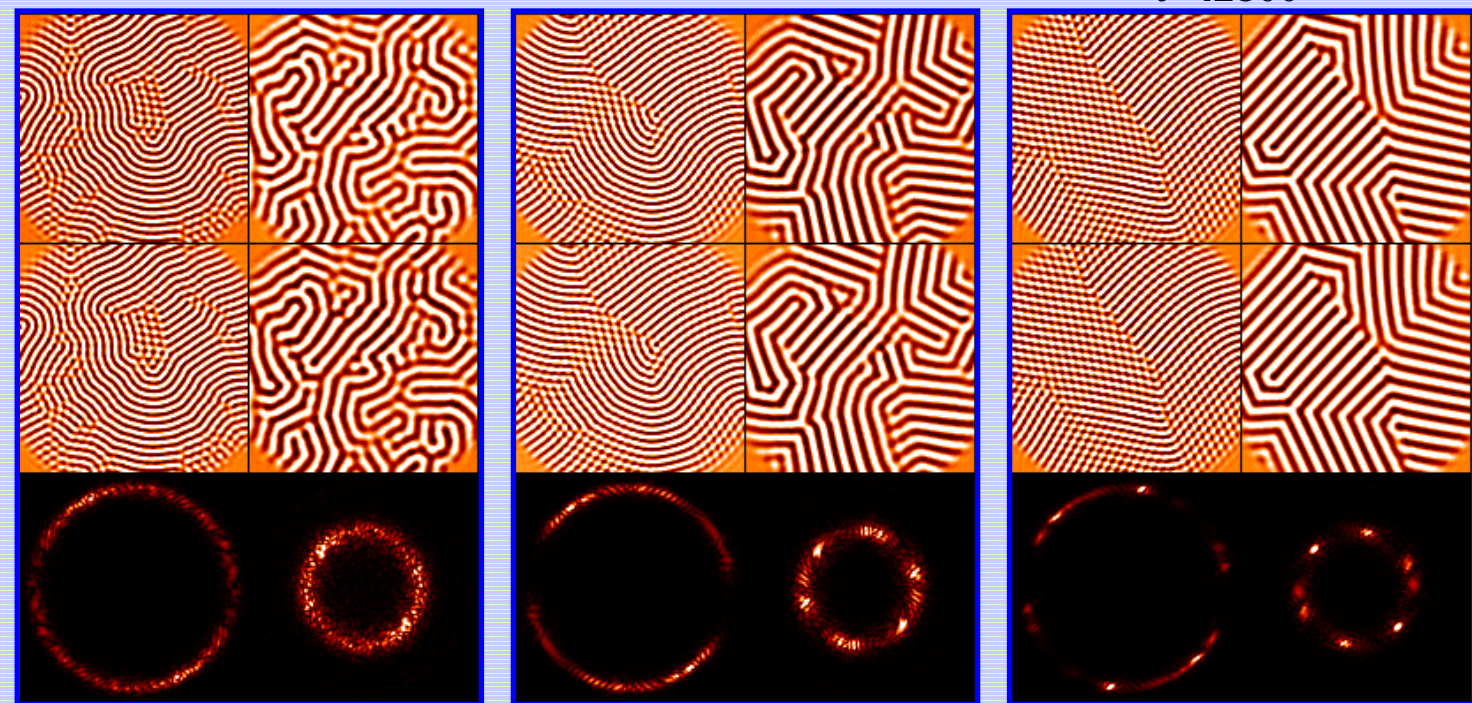


$|A_x| / FF$



$R_e(A_x) / R_e(A_y)$ →

$I_m(A_y) / I_m(A_x)$ →



Circular Polarized Intensity Patterns

Circular polarized states: $A_{\pm} = (A_X \pm i A_Y) / \sqrt{2}$

$$A_+ = (2i / \sqrt{2}) \text{Im} (A_X)$$

$$A_- = (2 / \sqrt{2}) \text{Re} (A_X)$$



•CP Intensity patterns with different wavelength

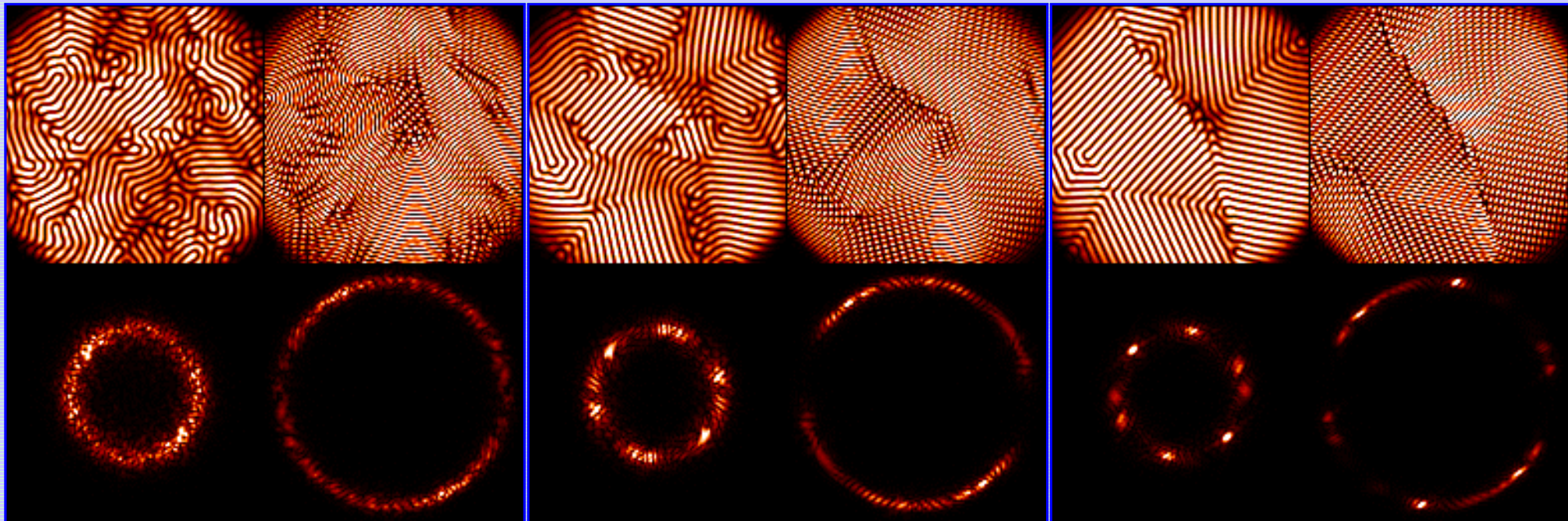
t=900

t=7800

t=42800

$|A_+|$

$|A_-|$

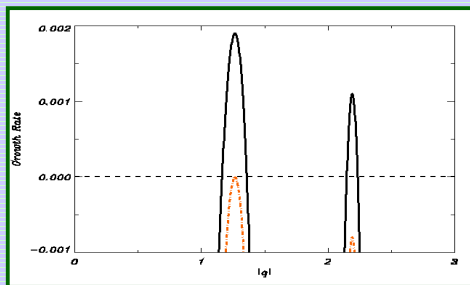


$FF(A_+)$

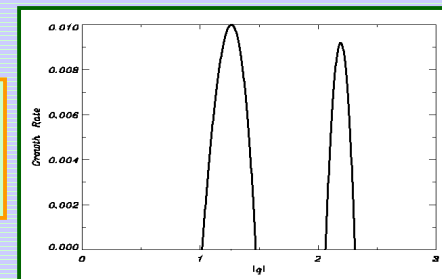
$FF(A_-)$

Type-II OPO, $\Delta_e < 0$, $\Delta_x \neq \Delta_y$, $\gamma_x \neq \gamma_y$, $\alpha_x \neq \alpha_y$

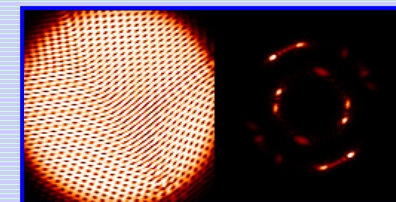
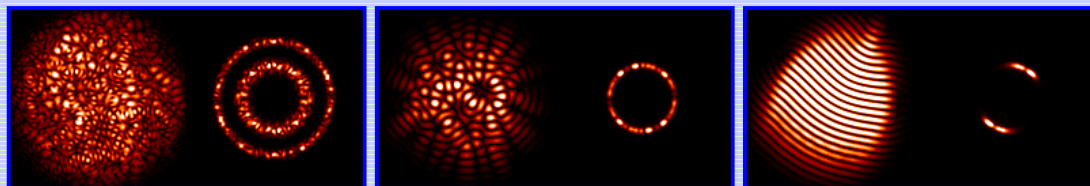
Near threshold



Far from threshold

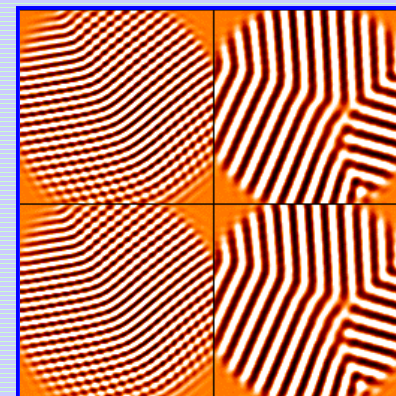
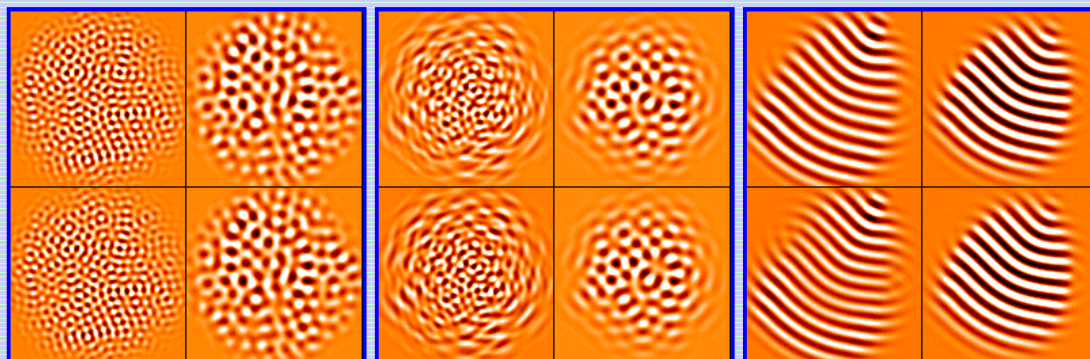


$|A_x| / FF$

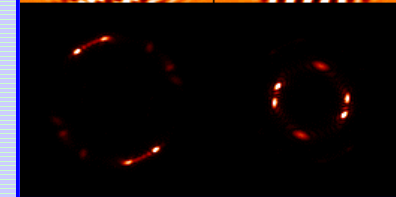
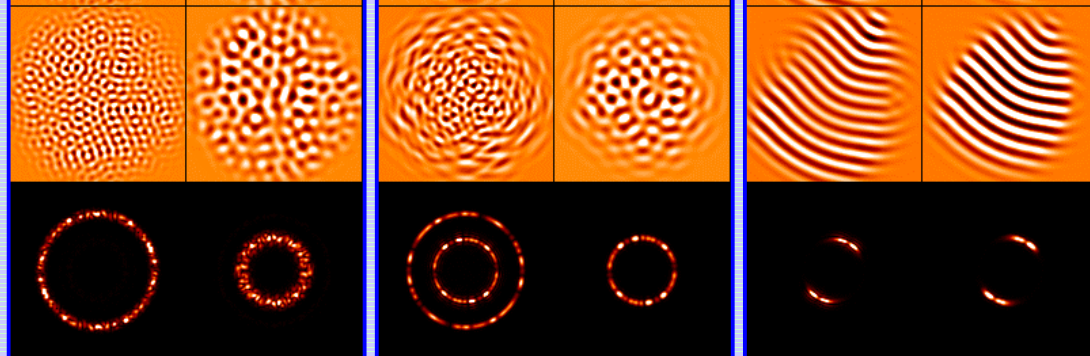


TIME →

$R_e(A_x) / R_e(A_y)$



$I_m(A_y) / I_m(A_x)$



t=100

t=3200

t=17600

t=15600

Amplitude Equations

Close to the instability threshold: Amplitude equations for the critical modes

$$\begin{aligned} \bullet A_1 &= A_x(q_1) \\ \bullet A_2 &= A_y(q_2) \end{aligned}$$

$$\begin{aligned} \partial_t A_1 &= \mu_1 A_1 - 4K_0^2 A_1 (|A_1|^2 + |A_2|^2) - \frac{2K_0^2}{1 + 4i\alpha' q_1^2} A_1 |A_1|^2 \\ \partial_t A_2 &= \mu_2 A_2 - 4K_0^2 A_2 (|A_1|^2 + |A_2|^2) - \frac{2K_0^2}{1 + 4i\alpha' q_2^2} A_2 |A_2|^2 \end{aligned}$$

$$\mu_{1,2} = \mu_{1,2}(F)$$

$$\eta_2(c', \Delta', \alpha', \mathbf{q}_{1,2})$$

• For F such that $\mu_1 > \mu_2 / (1 + \eta_2)$

Stable steady state

$$\begin{aligned} |A_1|^2 &\neq 0 \\ |A_2|^2 &= 0 \end{aligned}$$

Only patterns with wavenumbers q_1 are stable (near threshold)

• For F such that $\mu_1 < \mu_2 / (1 + \eta_2)$

Stable steady state

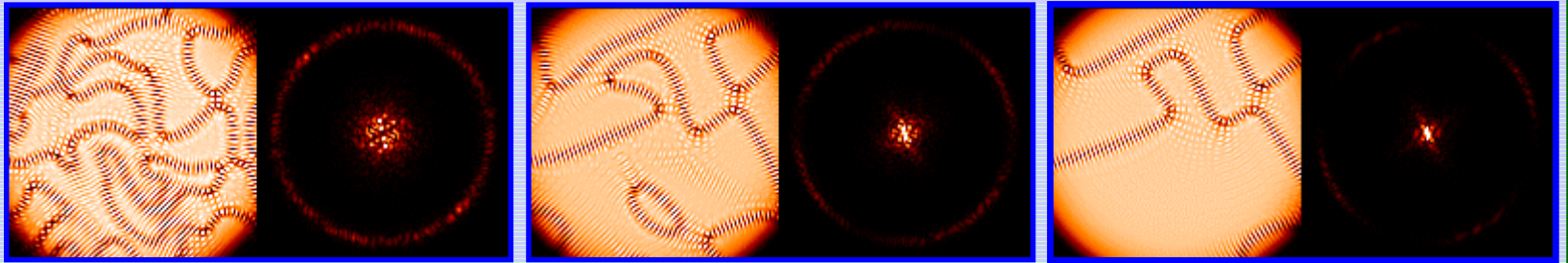
$$|A_1|^2 \neq |A_2|^2 \neq 0$$

Patterns with both wavenumbers $q_{1,2}$ are stable (far from threshold)

Type-II OPO, $\Delta_e < 0$, $q = 0$ mode selected

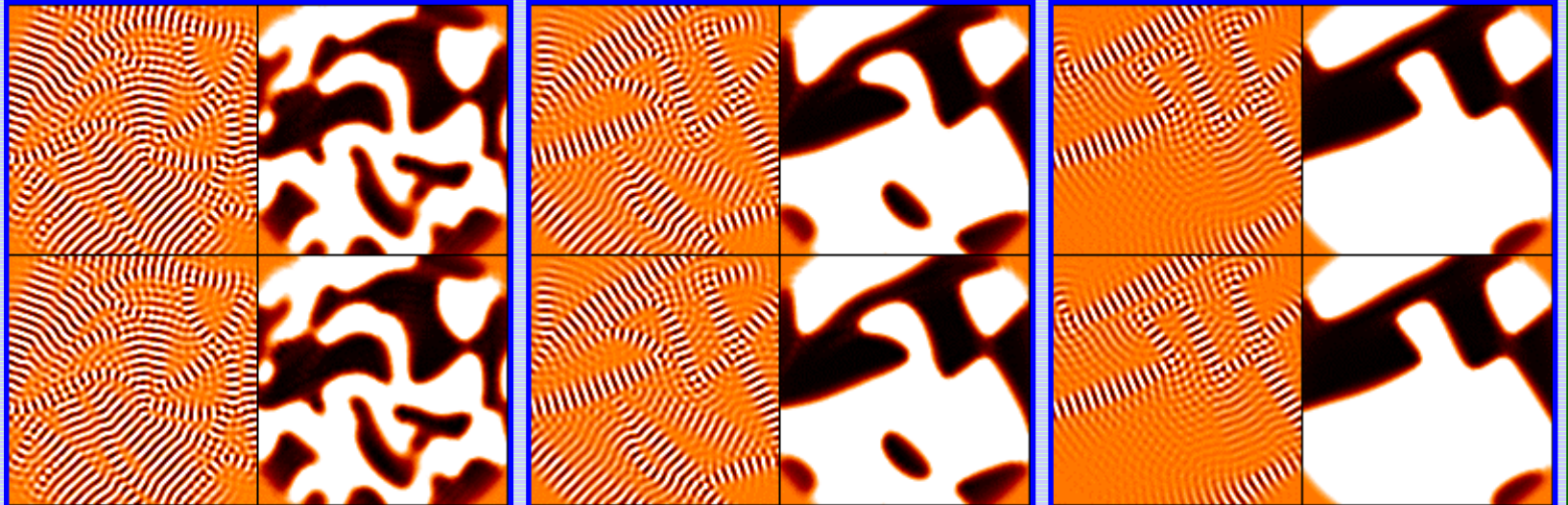
Symmetric coefficients

$|A_x| / FF$



$R_e(A_x) /$
 $R_e(A_y) \rightarrow$

$I_m(A_y) /$
 $I_m(A_x) \rightarrow$



t=400

t=3200

t=6000

SUMMARY

Birefringent / Dichroic coupling breaks relative phase invariance in Type II OPO : PHASE LOCKED states

$\Delta_e > 0$:

Inside the phase-locking regime

- Phase polarization domain walls
- Bloch Ising transition controlled by polarization coupling
- Core of the Bloch wall of orthogonal linear polarization
- Point defects on Bloch walls at points where chirality changes sign

Outside the phase-locking regime

- Oscillatory Bloch Domain Walls

$\Delta_e < 0$: •Pattern formation

- Standing waves for A_x and A_y
- Two competing modes in each linear polarization component
- Nonlinear mode selection: - mode coexistence far from threshold
- one mode selection close to threshold for asymmetric coefficients