

Simple models for scaling in phylogenetic trees

Emilio Hernández-García¹, Murat Tuğrul¹, E. Alejandro Herrada¹,
Víctor M. Eguíluz¹, Konstantin Klemm^{1,2}

¹IFISC (Instituto de Física Interdisciplinar y Sistemas Complejos), UIB-CSIC
Campus Universitat de les Illes Balears, E-07122 Palma de Mallorca, Spain

²Bioinformatics, Department of Computer Science, University Leipzig
Härtelstr. 16-18, 04107 Leipzig, Germany

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Abstract

Many processes and models produce trees with depth scaling logarithmically with the number of leaves. Phylogenetic trees, describing the evolutionary relationships between biological species, are examples of trees for which such scaling is not observed. With this motivation, we analyze numerically two branching models leading to non-logarithmic depth scaling. For the first one, Ford's alpha model, power-law scaling in the depth was established analytically. Our numerical results illustrate that the asymptotic regime is approached only at very large tree sizes. A second model, the *activity* model, is introduced here. We show analytically and numerically that its depth also displays power-law scaling at a critical parameter value.

1 Phylogenetic branching and models

Although most modern studies on complex networks [Albert & Barabási, 2002; Boccaletti et al., 2006] consider situations in which nodes are connected by multiple paths, the case of *trees*, i.e. graphs without closed cycles, is relevant to describe many natural and artificial systems. Branching in real trees [Stevens, 1974], in blood vessels [West et al., 1997], in river networks [Rodriguez-Iturbe & Rinaldo, 1997] or in computer file systems [Klemm et al., 2005, 2006] produce complex tree patterns worth to be described and understood. Trees are the outcome of classifications algorithms [Jain & Dubes, 1988] and of branching processes [Harris, 1963] and they also arise when computing

community structure [Guimerà et al., 2003] or as a backbone (for example a minimum spanning tree) for more connected networks [Garlaschelli et al., 2003; Hernández-García et al., 2007].

Evolutionary processes leading to speciation are also summarized in phylogenetic trees [Cracraft & Donoghue, 2004]. In these trees the leaves represent living species and each internal node represents a branching event in which an ancestral species diversifies into daughter species. Their topology encodes information on evolutionary mechanisms which is beginning to be scrutinized [Burlando, 1990, 1993; Ford, 2006; Blum & François, 2006; Hernández-García et al., 2007; Herrada et al., 2008].

The earliest mathematical model of evolutionary branching was proposed by Yule [1925]. Apart from the distinction he introduced between genera and species diversification, the model is equivalent to the Equal Rates Markov (ERM) model [Harding, 1971; Cavalli-Sforza & Edwards, 1967]: starting from a single ancestral species, one among the tree leaves existing at the present time is chosen at random, bifurcating into two new leaves. Then this operation is repeated for a number of time steps or, equivalently, until the tree reaches a desired size. The topological characteristics of the trees so constructed are surprisingly robust, being shared by apparently different models such as the coalescent and others [Aldous, 2001]. Essentially what is needed is that different branches at a given time branch independently and with the same probabilities. When extinction is taken into account, the same topology is recovered when considering only the lineages surviving at the final time. One of the characteristics of this

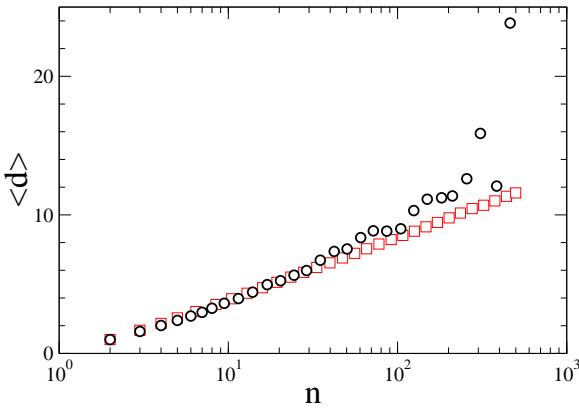


Figure 1: Mean depth $\langle d \rangle$ of trees in TreeBASE (circles) as a function of number of leaves n . Squares are obtained from computer simulations of the ERM model, behaving as Eq. (1) for large n . Depth in the real phylogenetic trees scales faster than the ERM behavior at large sizes. For both real phylogenies and model, depth values for each tree size are obtained by logarithmic binning of the depth of all trees and subtrees with that size.

type of branching is a distribution of subtree sizes A scaling at large sizes as A^{-2} , an outcome robustly observed in many natural and artificial systems, and in classification schemes, including taxonomies [Caldarelli et al., 2004; Capocci et al., 2008]. Another important characteristic is that the mean depth of the tree $\langle d \rangle$ (i.e. the average distance, measured in number of links, from leaves to the root) scales logarithmically with the number of leaves:

$$\langle d \rangle \sim \log n. \quad (1)$$

It is worth to note that these results apply not only to many random branching models, but also to the simple deterministic Cayley tree, in which all internal nodes at a given level split in a fixed number of daughter nodes.

In view of this generality it was surprising to find that the topology of observed phylogenies does not agree with any of these predictions [Herrada et al., 2008]. In fact, it was known since some time ago that real phylogenies are substantially more *unbalanced* than predicted by the ERM and similar models [Aldous, 2001; Blum & François, 2006]. This means that some lineages diversify much more than others, in a way that is statistically incompatible with the ERM predictions. Figure 1 compares data [Herrada et al., 2008] compiled from TreeBASE, a

public repository containing several thousands of empirical phylogenetic trees corresponding to virtually all kinds of organisms in Earth, with the predictions of the ERM model. For the phylogenetic trees at large sizes the mean depth scales with the number of leaves faster than the ERM behavior in Eq. (1).

The breakdown of the ERM behavior indicates that evolutionary branching should present correlations either in time or between the different branches. Mechanisms producing trees with non-ERM scaling for the depth have been identified, as for example the situation of critical branching [De Los Rios, 2001; Harris, 1963] or optimization of transport processes [Banavar et al., 1999]. In the phylogenetic context models of this type have been proposed [Aldous, 2001; Pinelis, 2003; Blum & François, 2006; Ford, 2006], although most of them lack a clear interpretation in biological terms.

In the following we present results for two branching models showing asymptotically non-ERM, i.e. non-logarithmic, scaling for the depth. Their study is motivated by the empirical results above from real phylogenetic trees, but we do not expect them to be good models for evolutionary process. Rather we study them because they pertain to the small set of available models with non-ERM scaling which are defined *dynamically* (i.e. by a set of rules that are applied to the present state of a growing tree to find the state at the next time step) rather than being characterized globally by statistical or optimization prescriptions. The first model we present, Ford's *alpha* model, is a simple example for which the non-trivial asymptotic scaling (of the power law type) has been analytically identified. We analyze it numerically to confirm this prediction and to display the behavior at finite sizes. We introduce later a new model, the *activity* model, which also leads to non-logarithmic depth scaling at a critical parameter value.

2 Ford's alpha model

Ford [2006] introduced a model for recursive tree formation: At a given step in the process the tree is a set of leaves connected by terminal links to internal nodes, which are themselves connected by internal edges until reaching the root (the root itself is considered to have a single edge joining to the first bifurcating internal

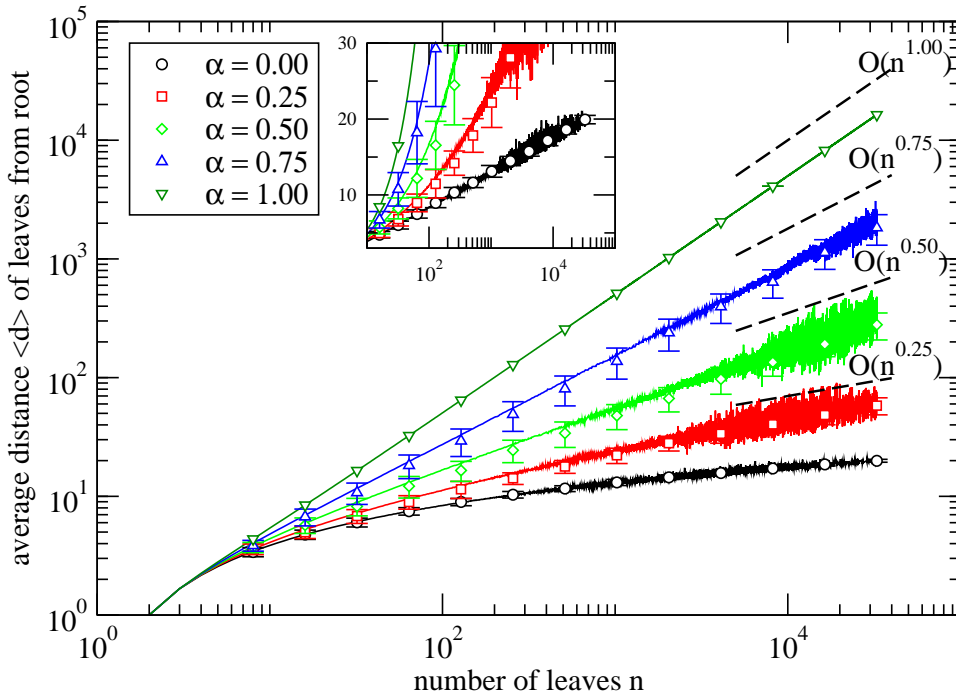


Figure 2: Depth statistics vs tree size for the alpha model. Symbols indicate the mean depth of leaves from root, averaged over the 100 trees generated for each size (2^k , $k = 3, 4, \dots, 15$), and the error bars are the corresponding standard deviations. The points in the rugged lines come from each subtree of all trees generated. The dashed segments indicate the analytic predictions [Ford, 2006] for the scaling at large n . The inset highlights the logarithmic scaling of the $\alpha = 0$ case.

node; with this convention a tree of n leaves has $n - 1$ internal edges). Then, a probability of branching proportional to $1 - \alpha$ is assigned to each leaf, and proportional to α to each internal edge. By normalization these probabilities are, respectively, $(1 - \alpha)/(n - \alpha)$, and $\alpha/(n - \alpha)$. When a leaf is selected for branching, it gives birth to a couple of new ones, as in the ERM model. But when choosing an internal edge, a new leaf branches from it by the insertion in the edge of a new internal node. For $\alpha = 0$ we have the standard ERM model. For $\alpha = 1$ the completely unbalanced comb tree, in which all leaves branch successively from a main branch, is generated. Intermediate topologies are obtained for $\alpha \in (0, 1)$.

By considering the effect of the addition of new leaves on the distances between root and other nodes, Ford [2006] derived exact recurrence relationships which, when written in terms of the average depth, lead to:

$$\langle d \rangle_{n+1} = \frac{n}{n - \alpha} \langle d \rangle_n + \frac{2n(1 - 2\alpha)}{(n + 1)(n - \alpha)}. \quad (2)$$

$\langle d \rangle_n$ is the mean depth of the leaves of a tree

with n leaves. By assuming a behavior $\langle d \rangle_n \sim n^\nu$ at large n , and expanding Eq. (2) in powers of $1/n$, we get $\nu = \alpha$, so that

$$\langle d \rangle_n \sim n^\alpha, \text{ if } 0 < \alpha \leq 1. \quad (3)$$

If $\alpha = 0$ the standard ERM behavior, Eq. (1), is recovered.

Figure 2 shows numerical results for the depth of trees generated with this model. Note that the predicted asymptotic behavior is attained, but only at very large tree sizes, in general much larger than available empirical phylogenies. As analytically demonstrated [Ford, 2006] depth statistics of subtrees of given size extracted from a large tree behave as data from trees of that size directly generated by the alpha model algorithm.

While the Ford model gives a simple mechanism for scaling in trees with a tunable exponent, the dynamical rule of posterior insertion of inner nodes is hard to justify in the context of evolution (although one can think on the modelling of errors arising in phylogenetic reconstruction methods when incorrectly assigning a splitting to a non-existing ancestral species).

This motivates the introduction of a new model described in the next section.

3 Activity model

In this section we show that tree shapes distinct from the ERM model may also result from a memory in terms of internal states of the nodes. The *activity* model proposed here is conceptually similar to the class of models suggested by Pinelis [2003]. However, the present model distinguishes only between active and inactive nodes and has a single parameter controlling the spread of activity.

Starting from a single disconnected node (the root), a binary tree is generated as follows. At each step, a leaf i of the tree is chosen and branched into two new leaves. Each of the two new leaves, independently of the other, is set active with probability p or inactive with probability $1 - p$. The branching leaf i is chosen at random from the set of active leaves if this set is non-empty. Otherwise, i is chosen at random from the set of all leaves. Figure 3 shows that for $p = 1/2$ the model generates trees with mean depth growing as the square root of tree size. For values of p below or above $1/2$, $\langle d \rangle$ seems to increase logarithmically with n .

Here we give a simplified argument to understand the observed exponent $1/2$ of the distance scaling with system size in the case $p = 1/2$. At the time the growing tree has n leaves in total, let $D_a(n)$ be the expected sum of distances of active leaves from the root, and $D_b(n)$ the analogous quantity for the inactive leaves. When a randomly chosen active leaf—at distance d_a from root—branches, the expected increase of $D_a(n)$ is

$$\begin{aligned} \Delta D_a(n) &= \\ p^2(d_a + 2) &+ 2p(1 - p) \cdot 1 + (1 - p)^2(-d_a) \\ &= (2p - 1)d_a + 2p. \end{aligned} \quad (4)$$

Here the three terms of the first sum are for the activation of two, one and zero of the new leaves. This expression is appropriate as far as the number of active nodes is not zero. Simultaneously, the expected change in $D_b(n)$ during the same event is

$$\begin{aligned} \Delta D_b(n) &= \\ p^2 \cdot 0 &+ 2p(1 - p)(d_a + 1) + (1 - p)^2 2(d_a + 1) \\ &= 2(1 - p)(d_a + 1). \end{aligned} \quad (5)$$

Averaging over the different active nodes which can be chosen amounts to replacing d_a in the above formulae by $\langle d_a \rangle_n$, the average depth of the active leaves in a tree of n leaves. Writing $D_i(n + 1) = D_i(n) + \Delta D_i(n)$, for $i = a, b$, one would get a closed system for the quantities $D_i(n)$ provided $\langle d_a \rangle_n$ is expressed in terms of them. This can be done by writing $\langle d_a \rangle_n = D_a(n)/a(n)$, where $a(n)$ is the expected number of active leaves in a tree of n leaves. This expected value is used here as an approximation to the actual number of active leaves.

The recurrence equations for $D_i(n)$ are specially simple in the most interesting case $p = 1/2$, since the dependence in $\langle d_a \rangle_n$ disappears from one of the equations:

$$D_a(n + 1) = D_a(n) + 1 \quad (6)$$

$$D_b(n + 1) = D_b(n) + \langle d_a \rangle_n + 1. \quad (7)$$

The solution (with initial condition $D_a(1) = 0$) of Eq. (6) is simply:

$$D_a(n) = n - 1. \quad (8)$$

Since the probabilities of an increment or decrement (by one unit) of the number of active leaves are the same and time-independent for $p = 1/2$, the number of active nodes performs a symmetric random walk with a reflecting boundary at 0 (this last condition arises from the prescription of setting active one node when the number of active nodes has reached zero in the previous step). For such random walk the expected value of active leaves $a(n)$ increases as the square root of the number of steps. Since a new leaf is added at each time step, this leads to:

$$a(n) \sim n^{1/2}. \quad (9)$$

Combining (8) and (9) we obtain the average distance of active nodes from root at large tree sizes:

$$\langle d_a \rangle_n \approx \frac{D_a(n)}{a(n)} \sim n^{1/2}. \quad (10)$$

Now we can plug this result into Eq. (7), which can be solved recursively:

$$D_b(n) = D_b(1) + \sum_{t=1}^{n-1} (\langle d_a \rangle_t + 1) \sim \sum_{t=1}^{n-1} t^{1/2} \sim n^{3/2}. \quad (11)$$

The totally averaged depth $\langle d \rangle_n$, which counts both the active and the inactive leaves, is

$$\langle d \rangle_n = \frac{D_a(n) + D_b(n)}{n} \sim \frac{n^{1/2} + n^{3/2}}{n} \sim n^{1/2}, \quad (12)$$

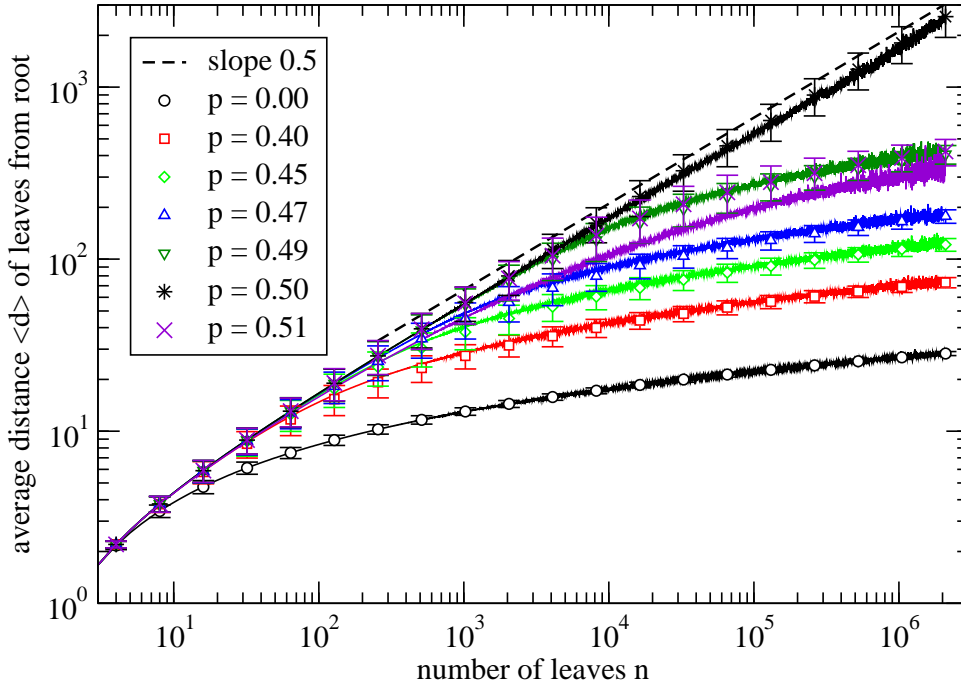


Figure 3: Depth versus size for the activity model for various values of the activation probability p . Data points displayed by symbols give the average distance of leaves with respect to the root. Error bars give the standard deviation taken over different realizations (1000 trees per data point). Data in the rugged curves are for all subtrees of trees with size $2^{21} = 2097152$. The dashed line has slope $1/2$, corresponding to the scaling of the $p = 0.5$ curve, as discussed in the text.

which explains the asymptotic behavior observed in Fig. 3 for $p = 1/2$.

We note that the growth dynamics presented here may be mapped to a branching process [Harris, 1963], with the difference that here the death (inactivation) of a node does not lead to its removal from the tree. The special case $p = 1/2$ corresponds to a critical branching process.

4 Discussion

We have presented and studied two simple models which lead to non-logarithmic scaling of the tree depth. In contrast with many of the available models having this behavior [Banavar et al., 1999; Aldous, 2001; Blum & François, 2006; Ford, 2006] they are formulated as dynamical models involving *growing trees*, so that rules are given to obtain the tree at the next time step from the present state. Their study has been motivated by data from phylogenetic branching, and they are interesting additions to our present understanding of complex networks and trees, but they can not be considered realistic representations of evolutionary

processes. The branching of internal edges in the Ford model has no obvious biological interpretation. The activity model puts the mechanisms of birth-death critical branching [Harris, 1963] within a framework of transitions between node internal states similar in spirit to the approach of Pinelis [2003]. The need to tune a parameter to attain the non-ERM critical behavior is however a limitation for its applicability. Much additional work is needed to identify the proper biological mechanisms behind evolutionary branching and adequate modelling of them.

A recent analysis of several evolutionary models including species competition [Stich & Manrubia, 2008] indicates that in these models correlations are finally destroyed by mutation processes and persist only for a finite correlation time. Thus sufficiently large trees would have a scaling behavior closer to the asymptotic ERM predictions. Since the largest phylogenies in databases such as TreeBASE have only some hundreds of leaves, it is possible that the observed imbalance and depth scaling is a finite-size regime. Nevertheless models going beyond

the ERM scaling are needed at least to explain this finite-size regime, and also to elucidate the true asymptotic scaling behavior. Here, we have also observed large finite-size transients in the alpha model of Sect. 2.

We finally note that since $\langle d \rangle$ is a measure of the *diameter of the tree*, and that for a tree the number of nodes is simply twice the number of leaves, power law scaling of the type $\langle d \rangle \sim n^\nu$ indicates that the tree can be thought as having a (fractal) dimension $d = 1/\nu$ [Eguíluz et al., 2003]. The logarithmic scaling in the ERM model is an example of the *small-world* behavior common to many network structures [Albert & Barabási, 2002], indicating an effective infinite dimensionality, whereas the power law scaling reveals a finite dimension for the tree. The alpha model produces trees with tunable dimension from 1 to ∞ , and the critical activity model gives two-dimensional trees.

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